

Quantum Sensors and Axion Dark Matter Searches

Wright Lab Quantum Sensing Workshop

Quantum Week at Yale

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U.S. DEPARTMENT OF
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QuantiSED



QISE-NET
QUANTUM INFORMATION SCIENCE
AND ENGINEERING NETWORK

Axions are well motivated

Axion Dark Matter

$a \leftrightarrow \gamma\gamma$ Parameter Space

Present day axion density

$$\Omega_a h^2 \approx 0.1 \left(\frac{10 \mu\text{eV}}{m_a} \right)^{7/6} \langle \theta_i^2 \rangle$$

Initial misalignment

Pre-Inflationary PQ Breaking

(f_a near GUT scale)

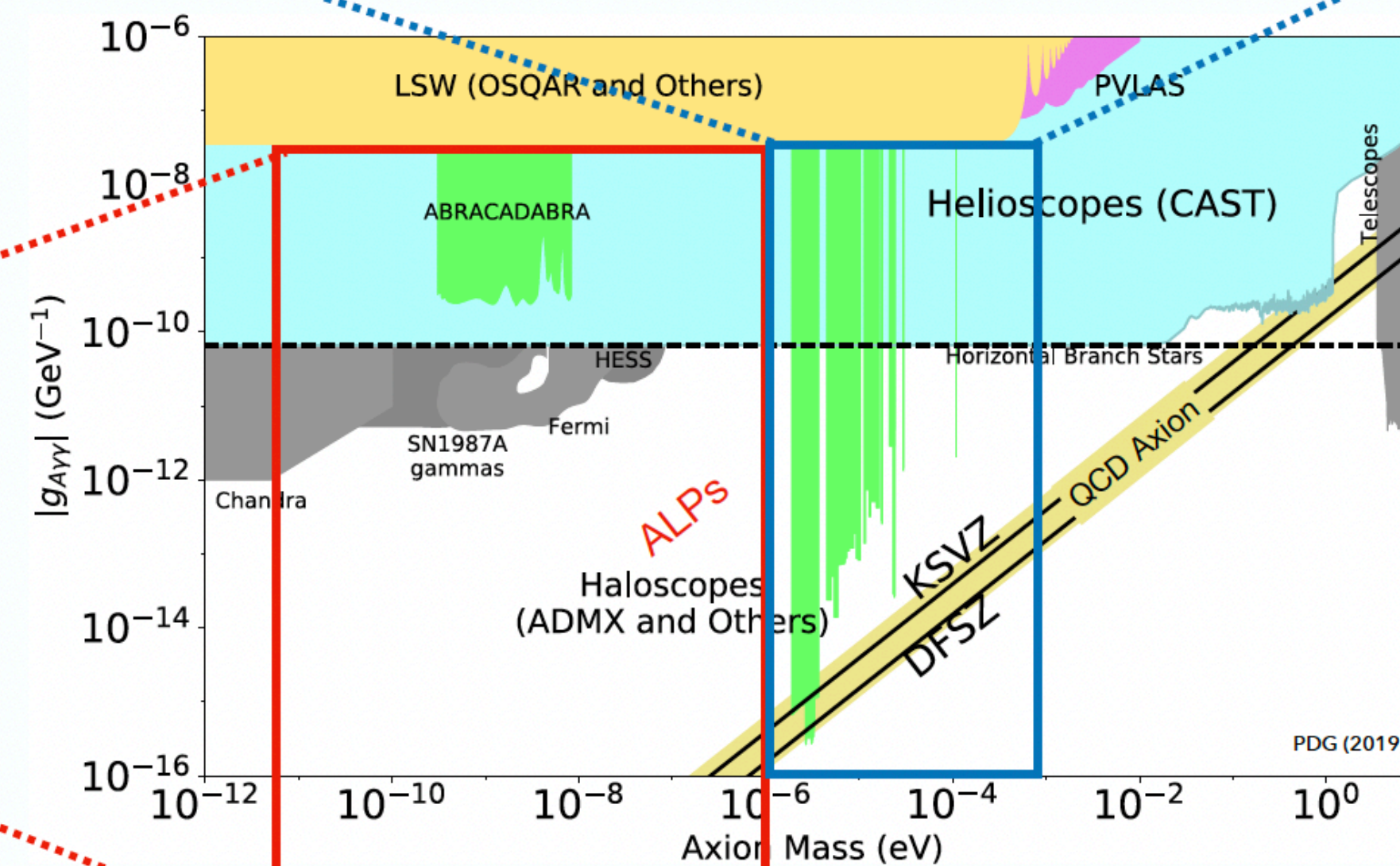
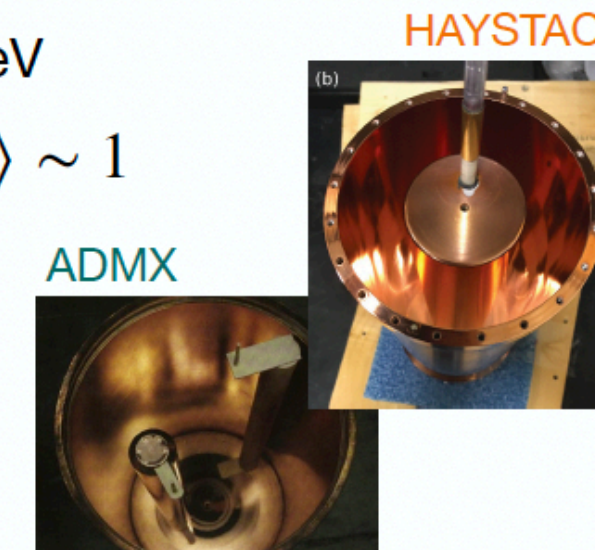
- Mass range $20 \text{ peV} \lesssim m_a \lesssim 1 \mu\text{eV}$
- Strong particle physics argument "GUT-scale" axion ($f_a \sim 10^{17} \text{ GeV}$)
- Small initial misalignment $\langle \theta_i^2 \rangle < 1$
- Long Compton wavelength regime (Magneto quasistatic regime)
- Lumped element detectors



Post Inflationary PQ Breaking

($f_a \sim 10^{12} \text{ GeV}$)

- Mass range $1 \mu\text{eV} \lesssim m_a \lesssim 1 \text{ meV}$
- Large initial misalignment $\langle \theta_i^2 \rangle \sim 1$
- Microwave Cavity regime
- ADMX, HAYSTAC, CAPP-8TB, QUAX-ay, ORGAN, others...



J. Ouellet

Aspen Winter Conference 2022

March 23, 2022

Axion searches at Yale



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Haystack  + Rydberg extension



Steve Lamoreaux



Reina Maruyama



Sid Cahn



James Nikkel



Danielle Speller
-> Johns Hopkins



Mike Jewell



Yuqi Zhu



Sumita Ghosh



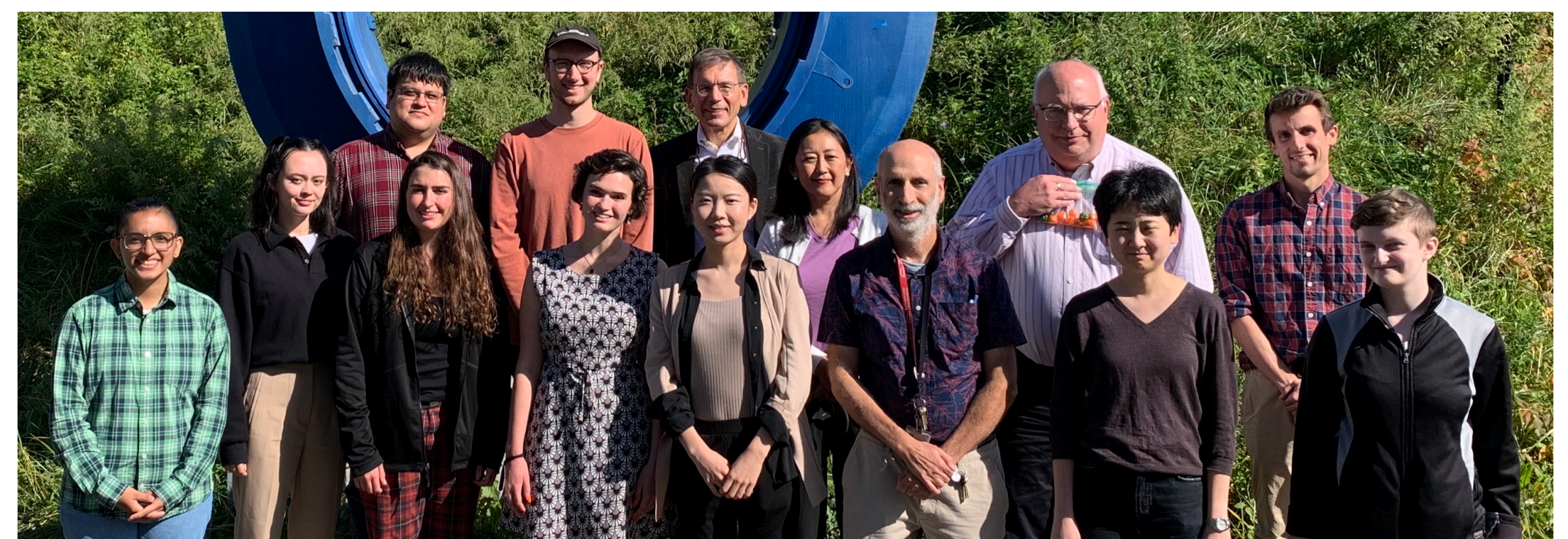
Xiran Bai



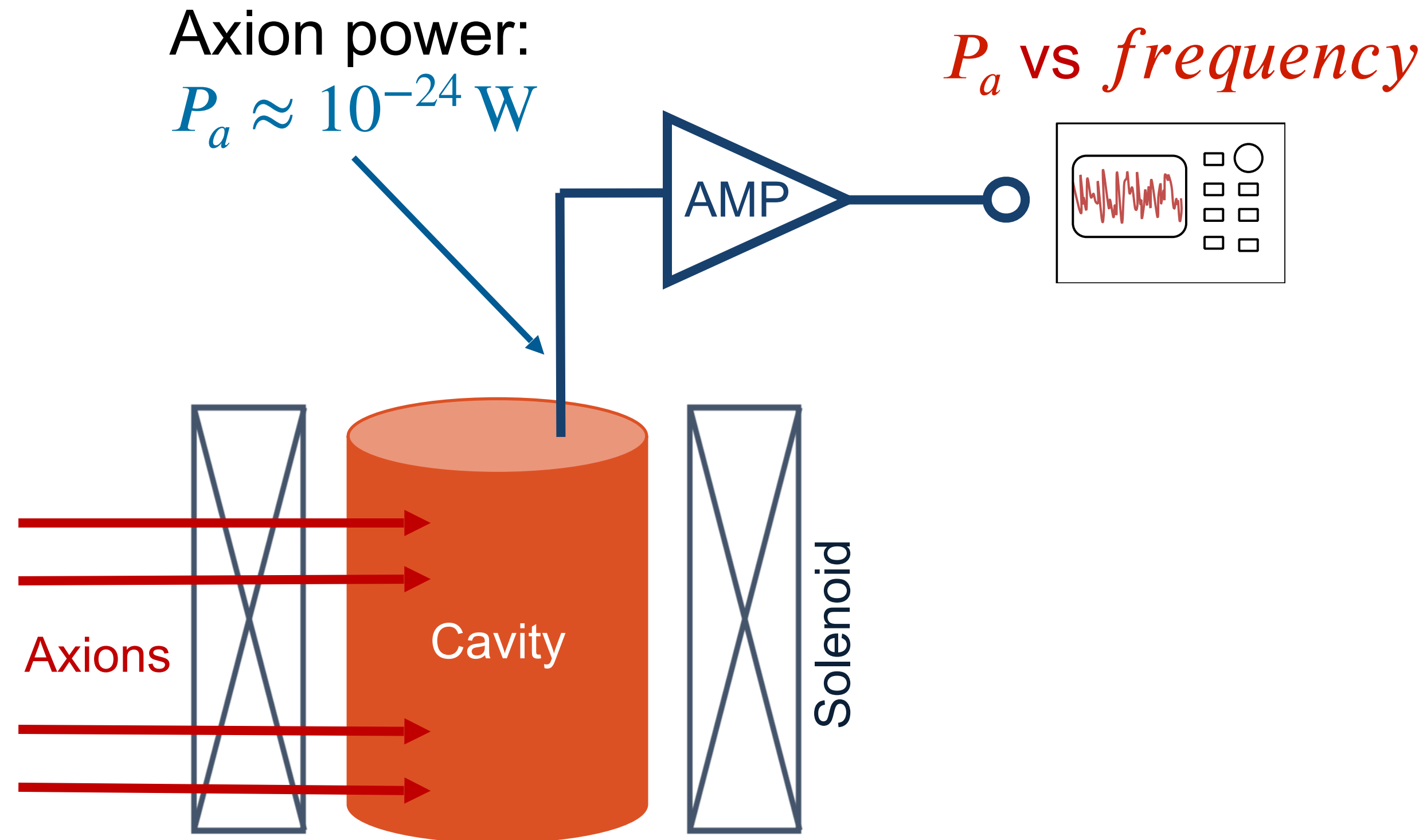
Eleanor Graham



Claire Laffan



Basic axion detection principle

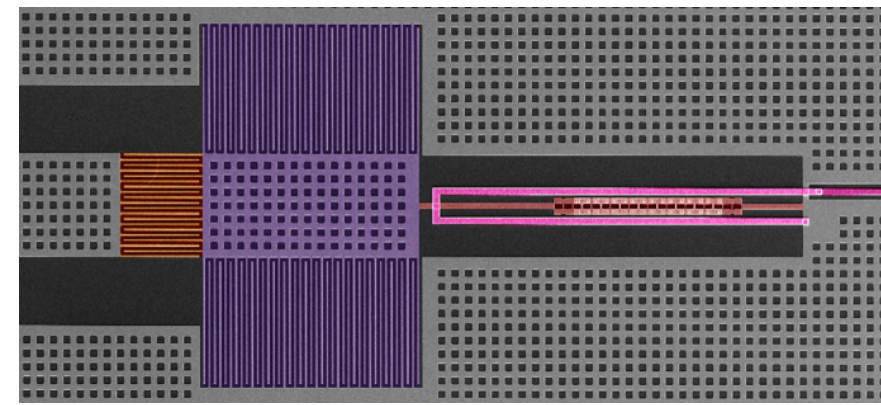


Interaction of interest: $\mathcal{L} \supset g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$

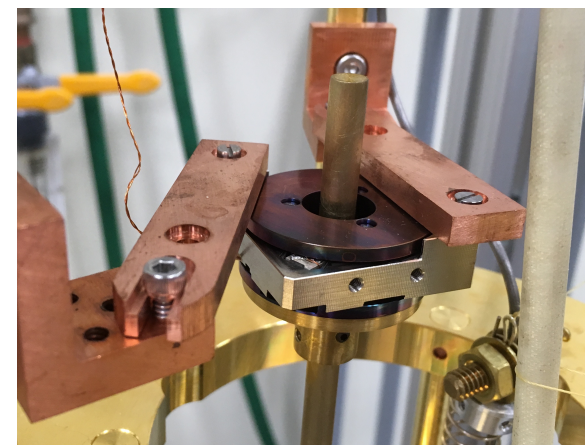
- High Q cavity: $Q = \frac{f_c}{\Delta f_c}$
- Low noise amplifier
- Tunable: $hf_a \approx m_a c^2$
- Large magnet: $B = 8 \text{ T}$
- Cryogenic: $T = 60 \text{ mK}$

Haloscope principle: P. Sikivie, *Phys. Rev. Lett.*, **51**, 1415 (1983)
 HAYSTAC detector: *Nucl. Instrum. Methods A* 854, 11 (2017)

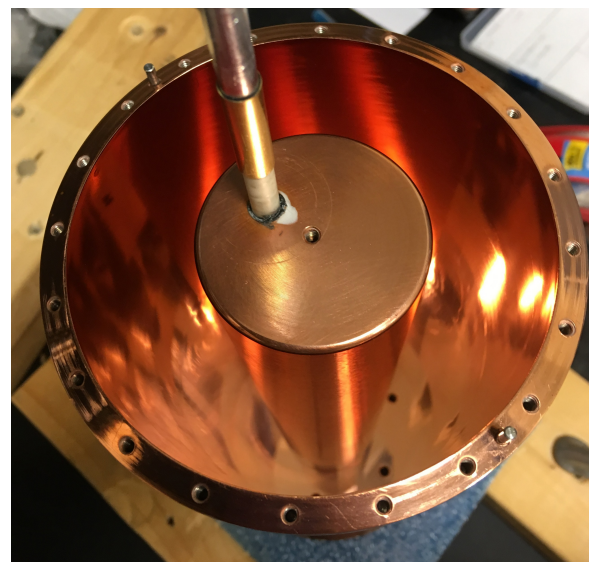
HAYSTAC Experiment



Quantum amplifiers



Piezoelectric tuning



Microwave cavity

Cryostat
operating temp
~60 mK

9 T magnet

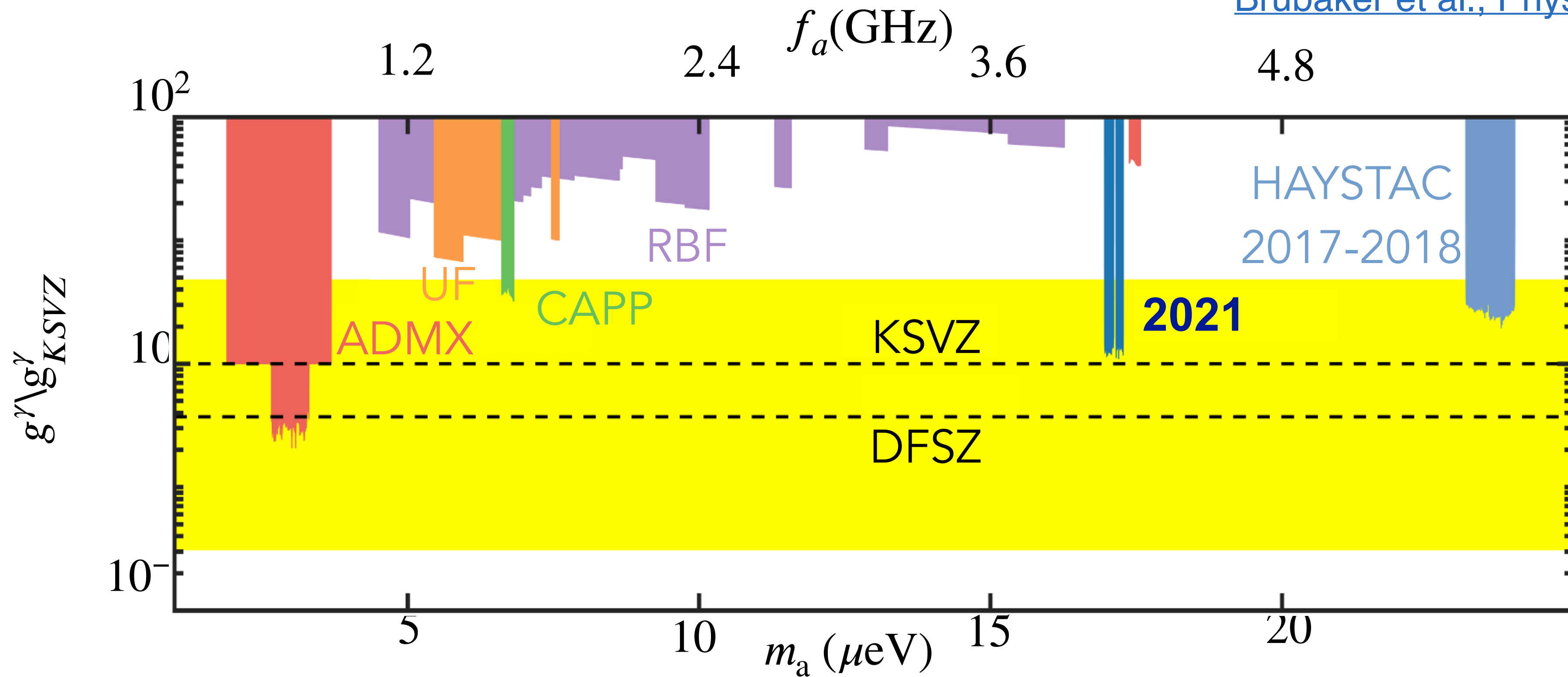


HAYSTAC results

[Backes et al., Nature, 590, 238–242 \(2021\)](#)

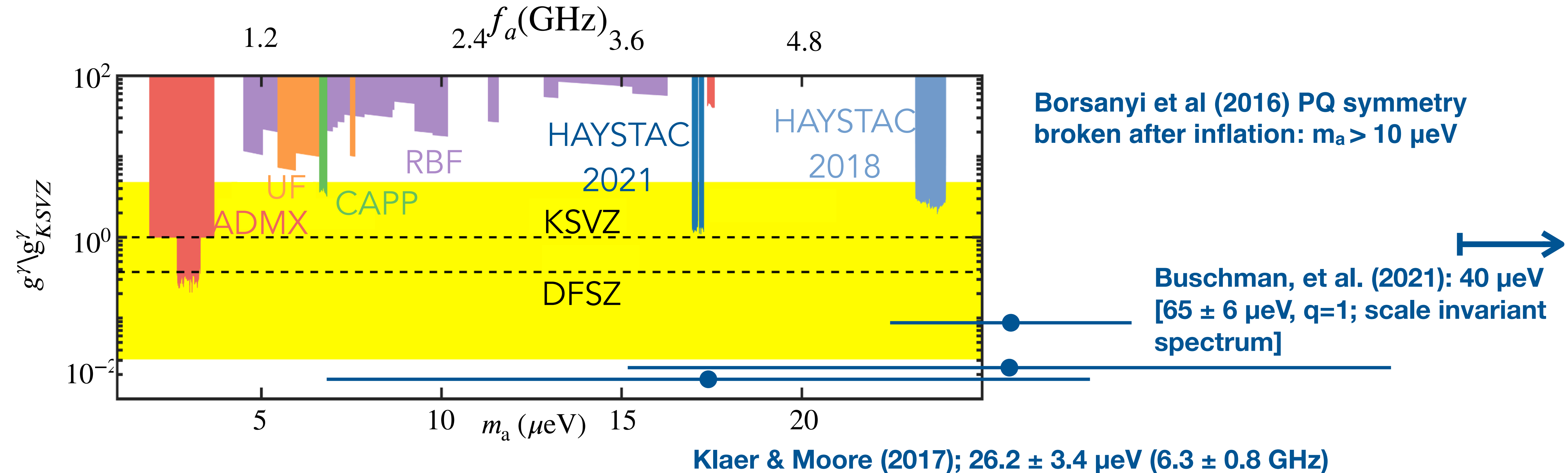
[Zhong et al., Phys. Rev. D 97 092001 \(2018\)](#)

[Brubaker et al., Phys. Rev. Lett. 118 061302 \(2017\)](#)



- dark matter exclusion enhanced by quantum squeezing
- sensitive axion search, dipping into KSVZ $> 10 \mu\text{eV}$
- New data 2020 – 2021, stay tuned

Axions well-motivated @ $m_a > 15 \mu\text{eV}$



Buschman, Foster & Safdi (2019):
 $25.2 \pm 11 \mu\text{eV}$ (6.1 \pm 2.7 GHz)
 $17.4 \pm 11 \mu\text{eV}$ (4.2 \pm 1.1 GHz)*

* In $\Omega_A \sim f_A^\alpha$, the best fit $\alpha = 1.24 \pm 0.04$
 Rather than analytical 1.187

Single photon detectors

- Single photon detectors have lower noise at higher frequencies Lamoreaux et al., 2013

Linear amplifier

- low-T Dicke radiometer eq.

$$P_l = h\nu(\bar{n}_{\text{BBR}} + 2 \times \frac{1}{2}) \sqrt{\frac{\Delta\nu_a}{\tau_{\text{int}}}}$$

($\bar{n}_{\text{BBR}} = \frac{1}{e^{h\nu/k_B T} - 1}$, 1/2: zero-point fluctuations)

- sensitive to amplitude and phase

Single-photon detector

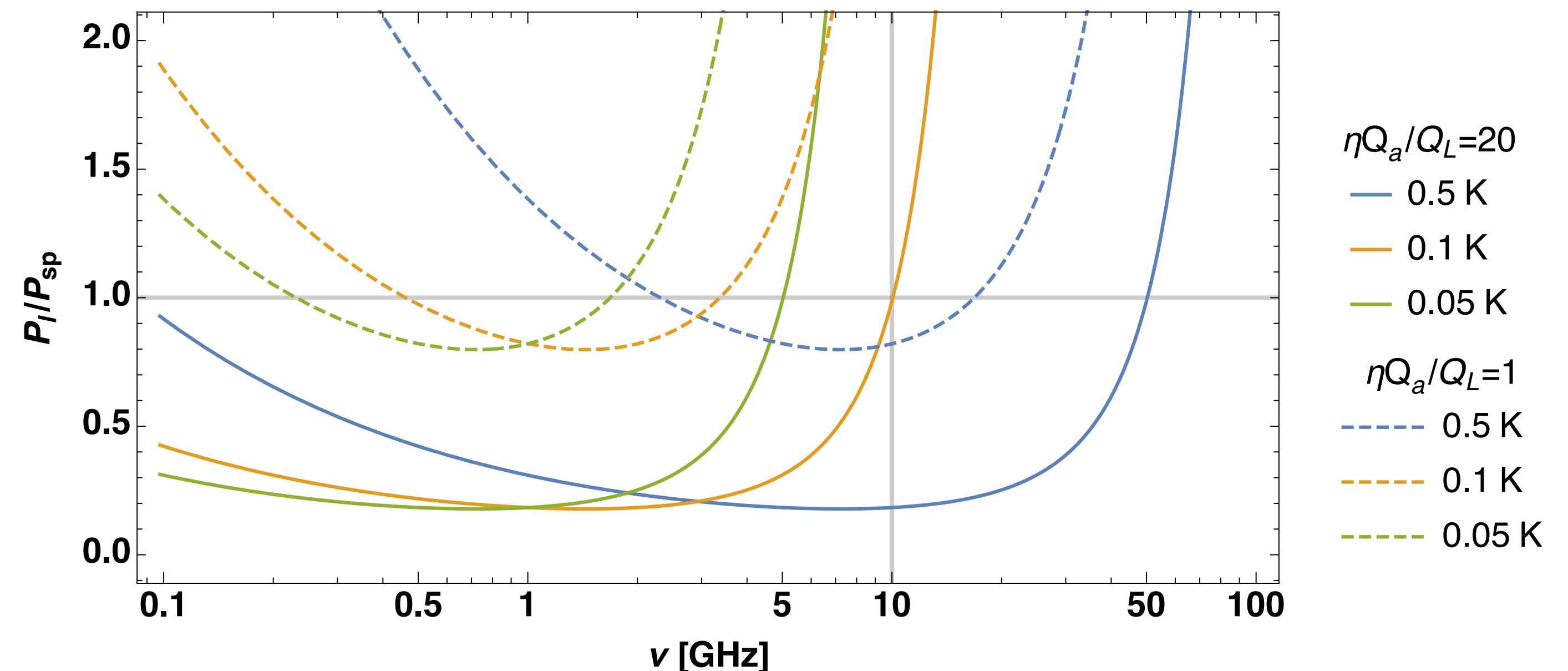
- shot noise due to BBR

$$P_{\text{sp}} = h\nu \sqrt{\frac{\eta \bar{n}_{\text{BBR}} \Gamma}{\tau_{\text{int}}}}$$

(η : efficiency, $\Gamma = 2\pi\nu/Q_L$ decay rate)

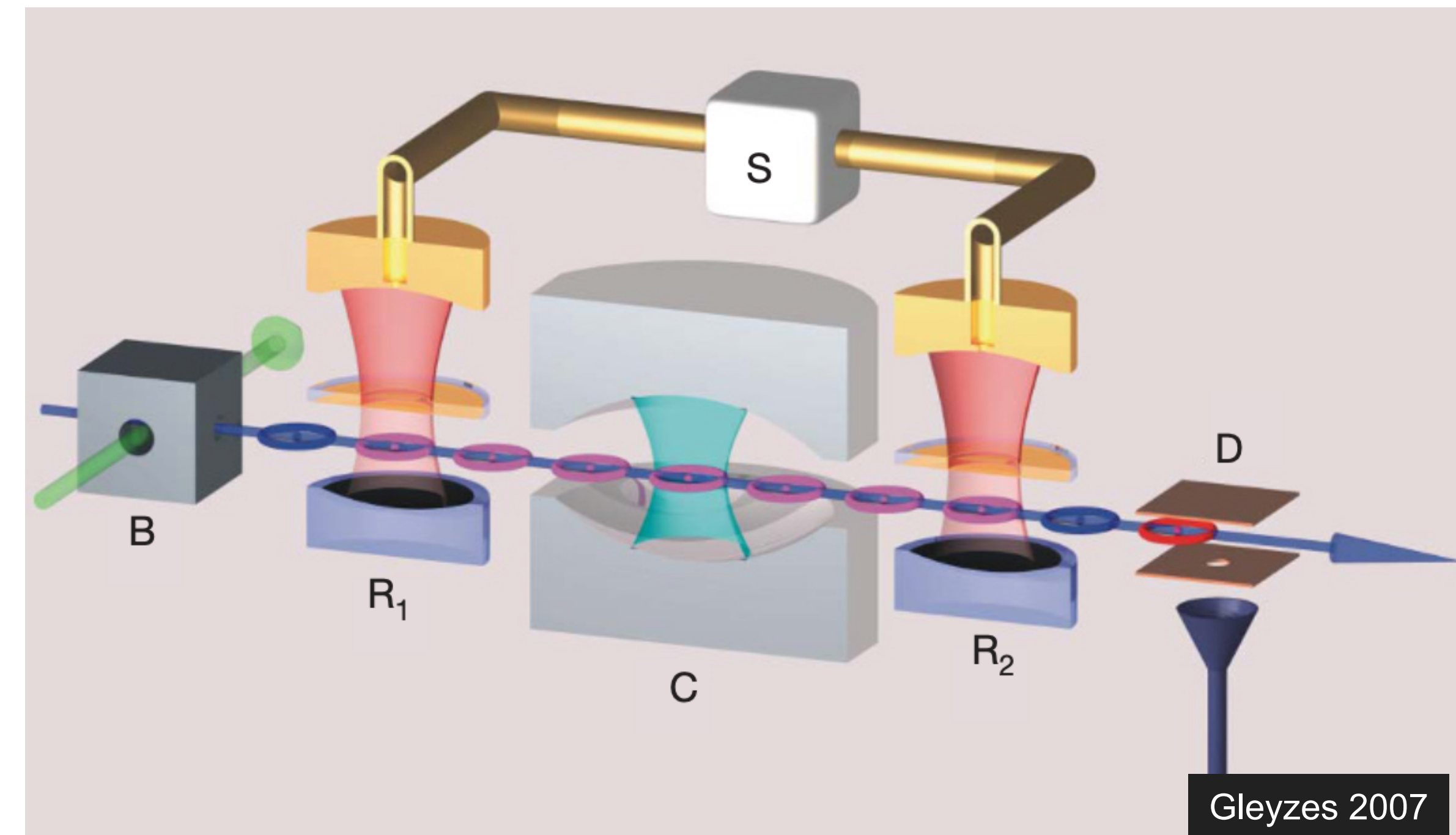
- only sensitive to amp.

$$\frac{P_l}{P_{\text{sp}}} = \frac{\bar{n}_{\text{BBR}} + 1}{\sqrt{\bar{n}_{\text{BBR}}}} \sqrt{\frac{Q_L}{2\pi\eta Q_a}}$$

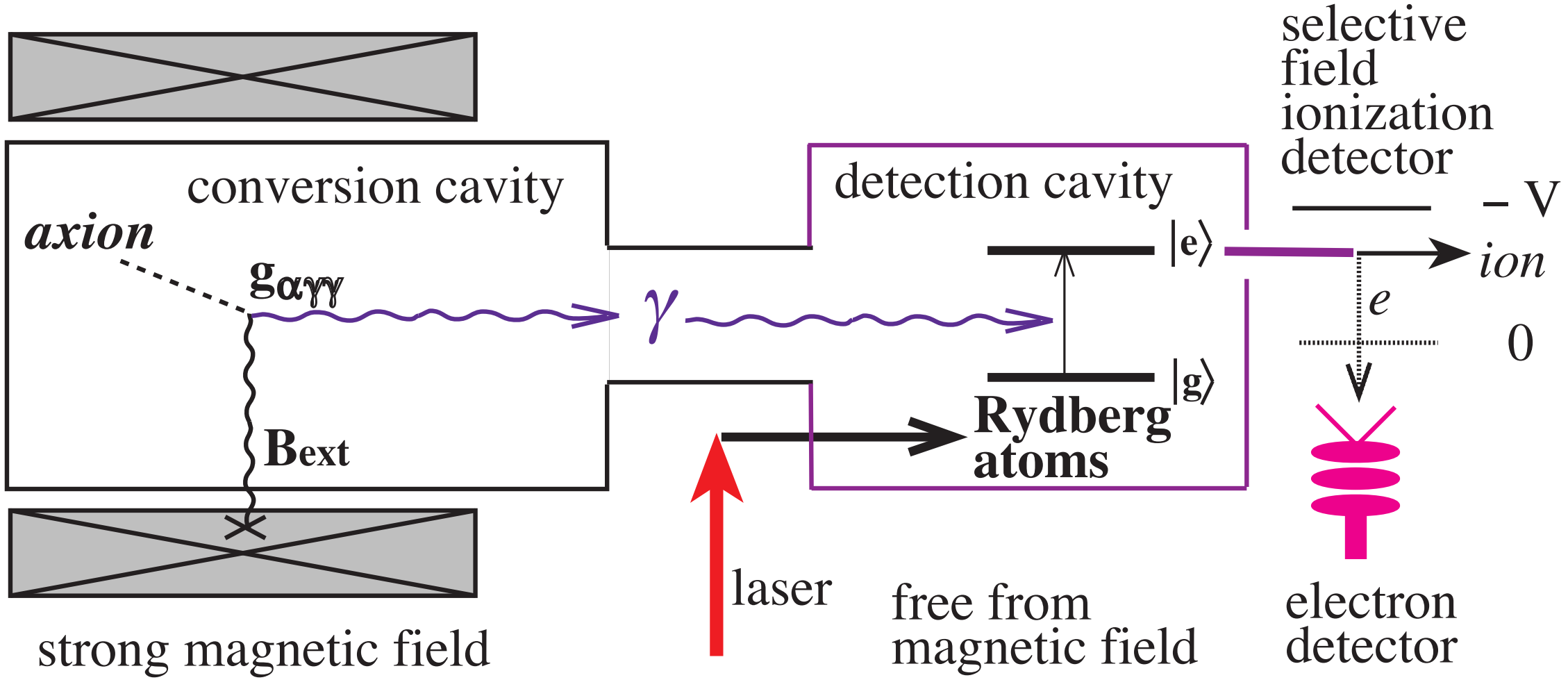


Rydberg atoms as microwave detectors

- Rydberg atoms:
 - Highly excited valence e^-
 - Couple strongly to 10 - 1000 GHz
- Applications:
 - Vapor-cell electrometry
 - e.g. Stuttgart; Sediacek et al., 2016
 - Single-atom cavity QED: ENS
 - **CARRACK axion search**

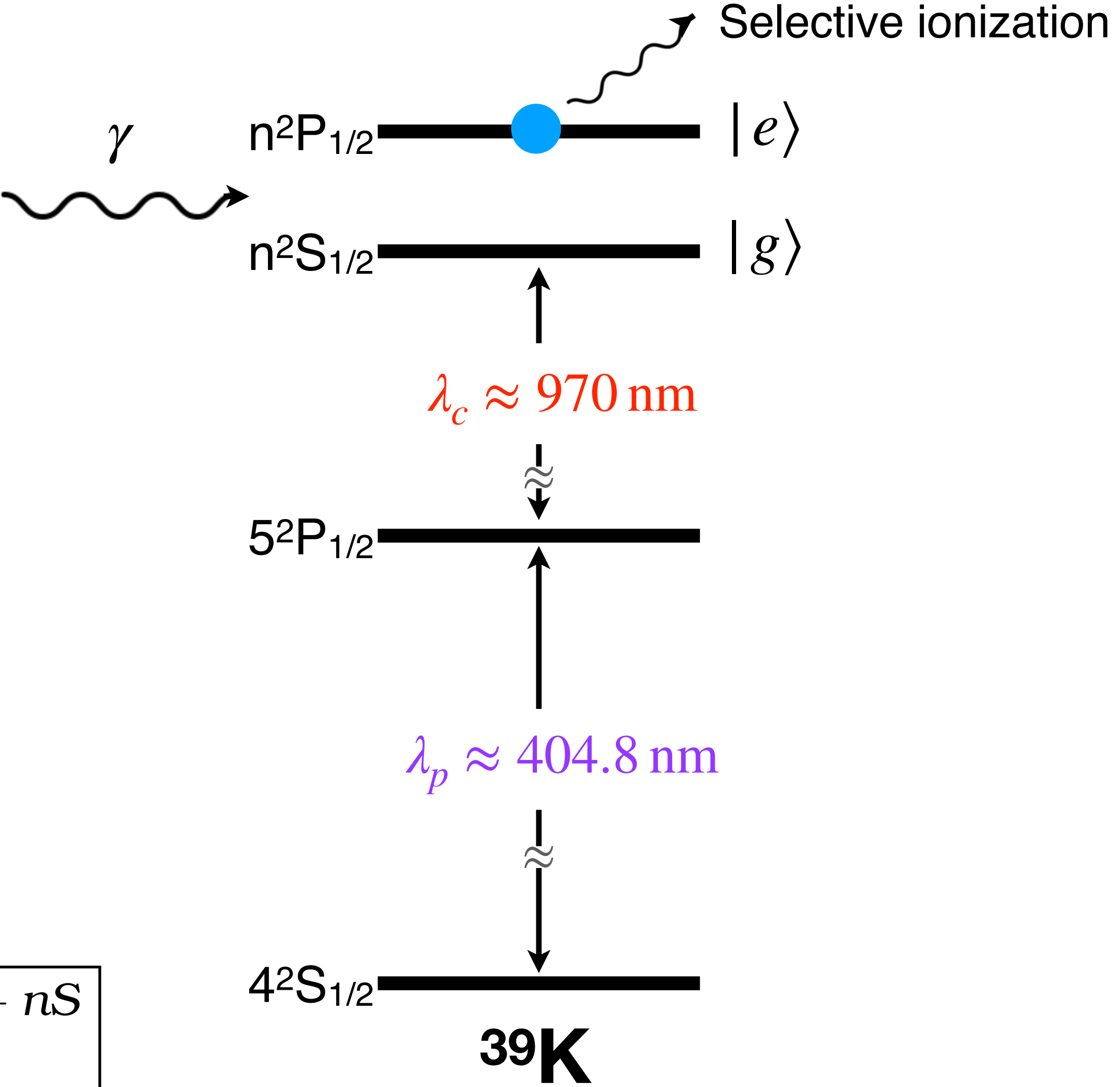


Basic principle



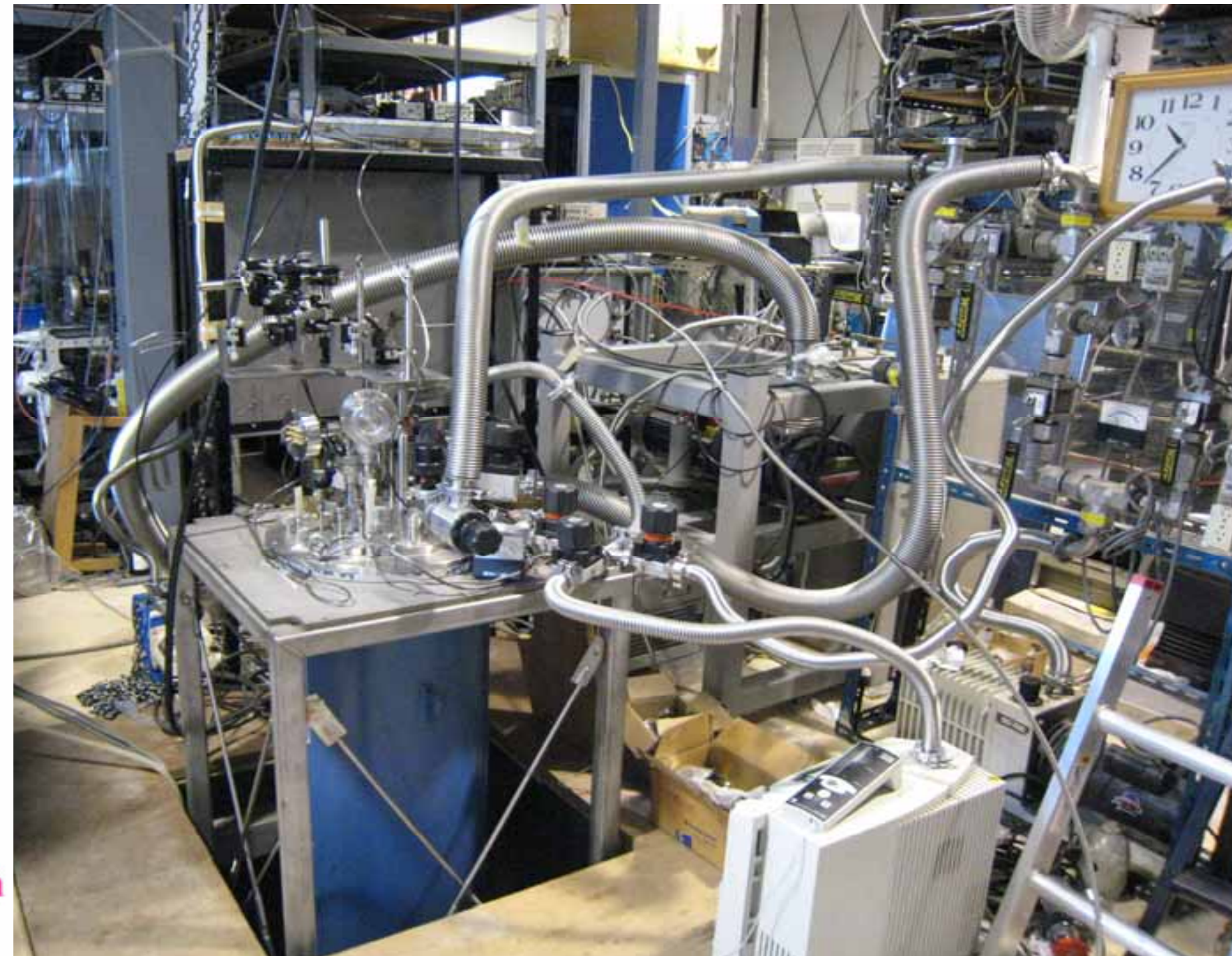
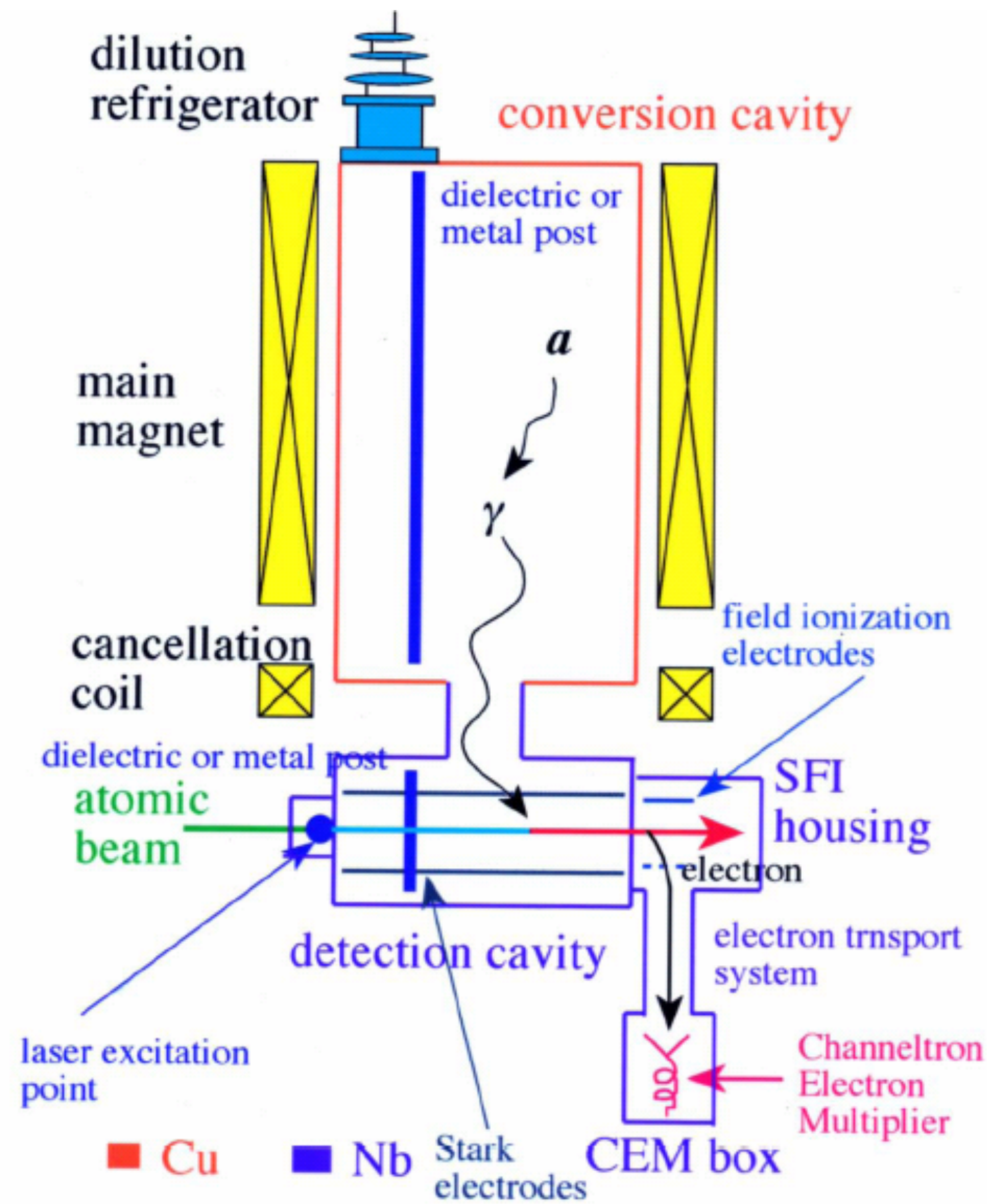
alternative: $4S - 4P_{3/2} - nS$
@767 nm, 457 nm

Rydberg atom

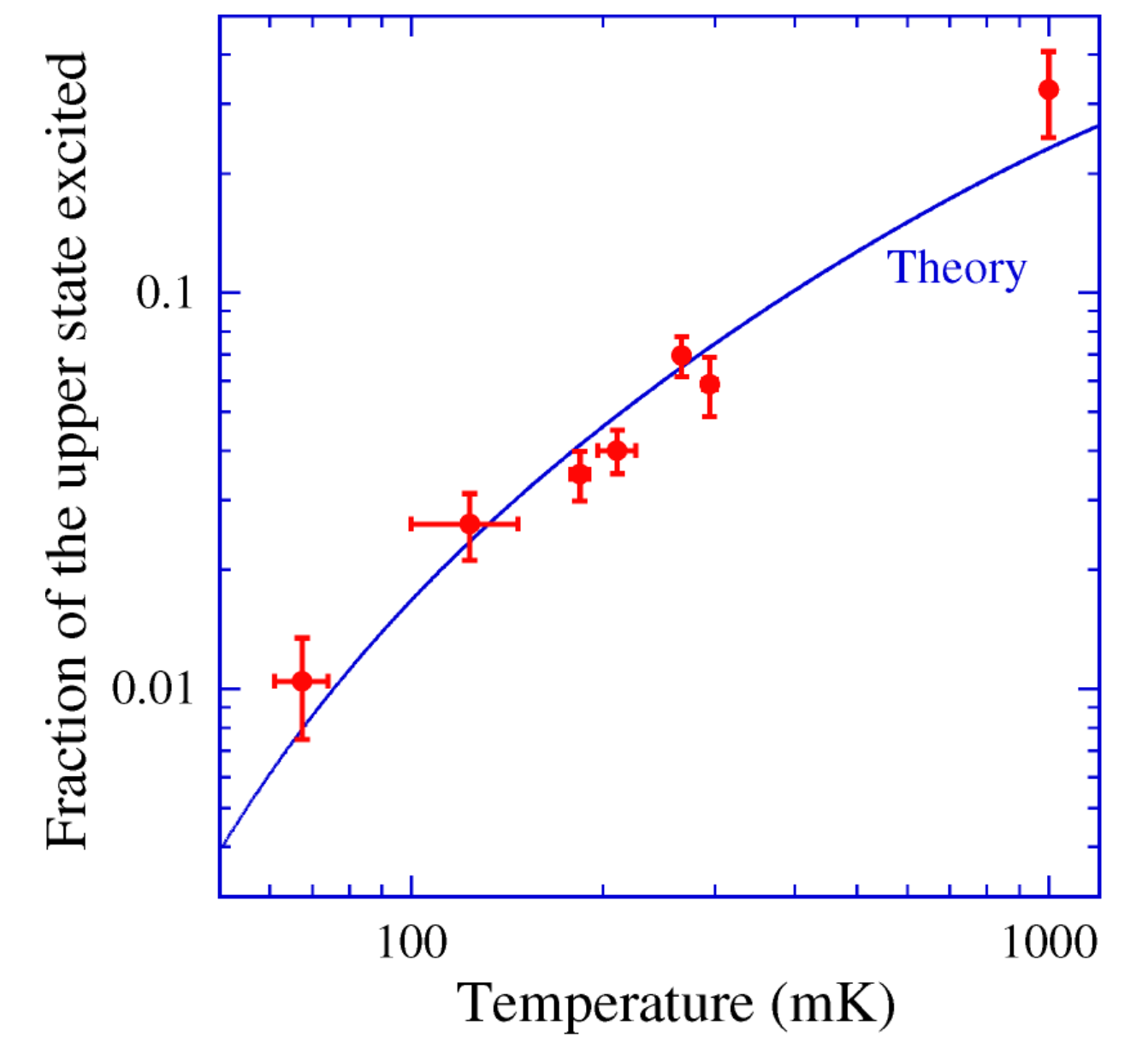


CARRACK

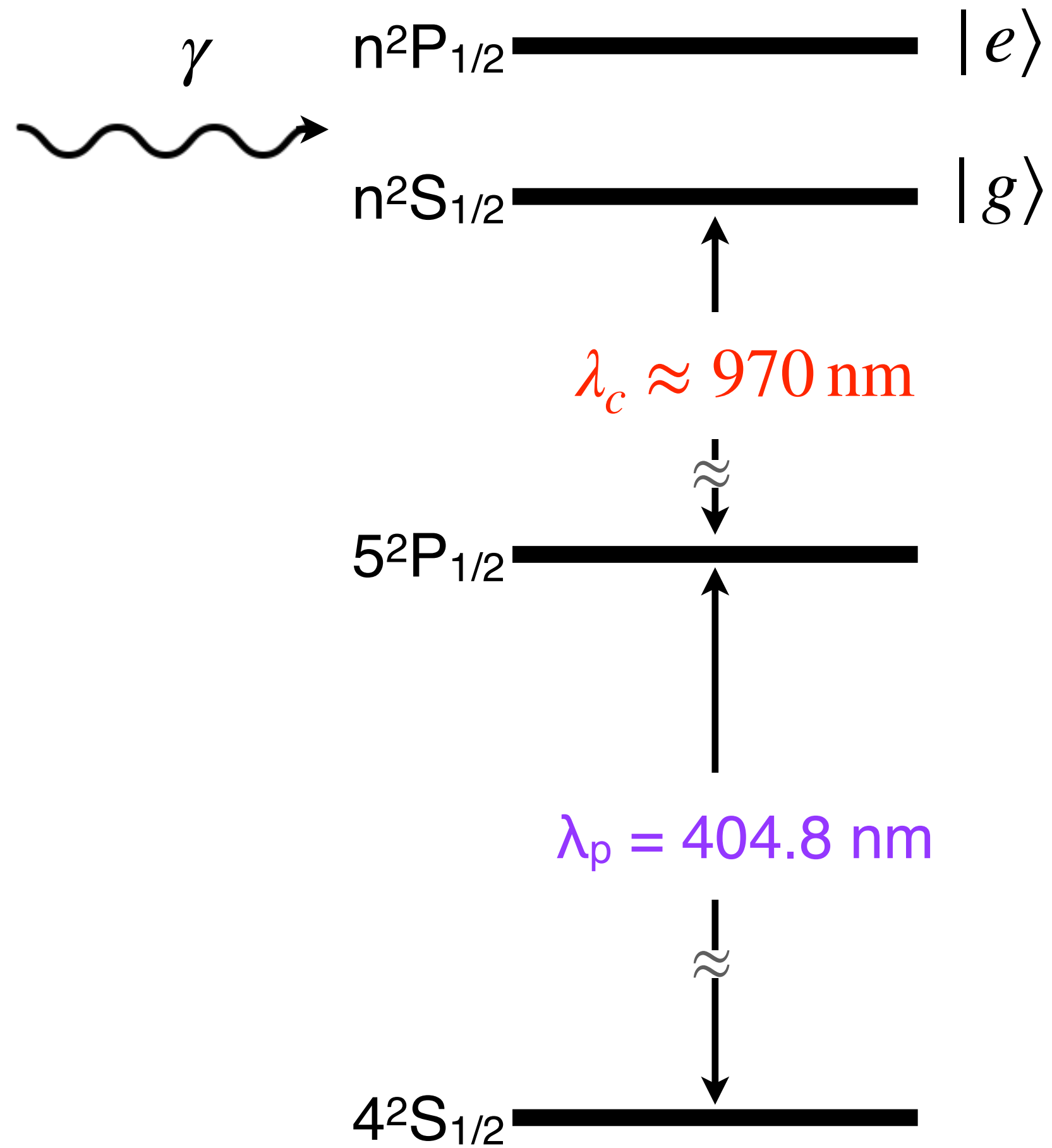
Imai, PANIC 2008
Tada, PLA 349 (2006) 488



Black body radiation



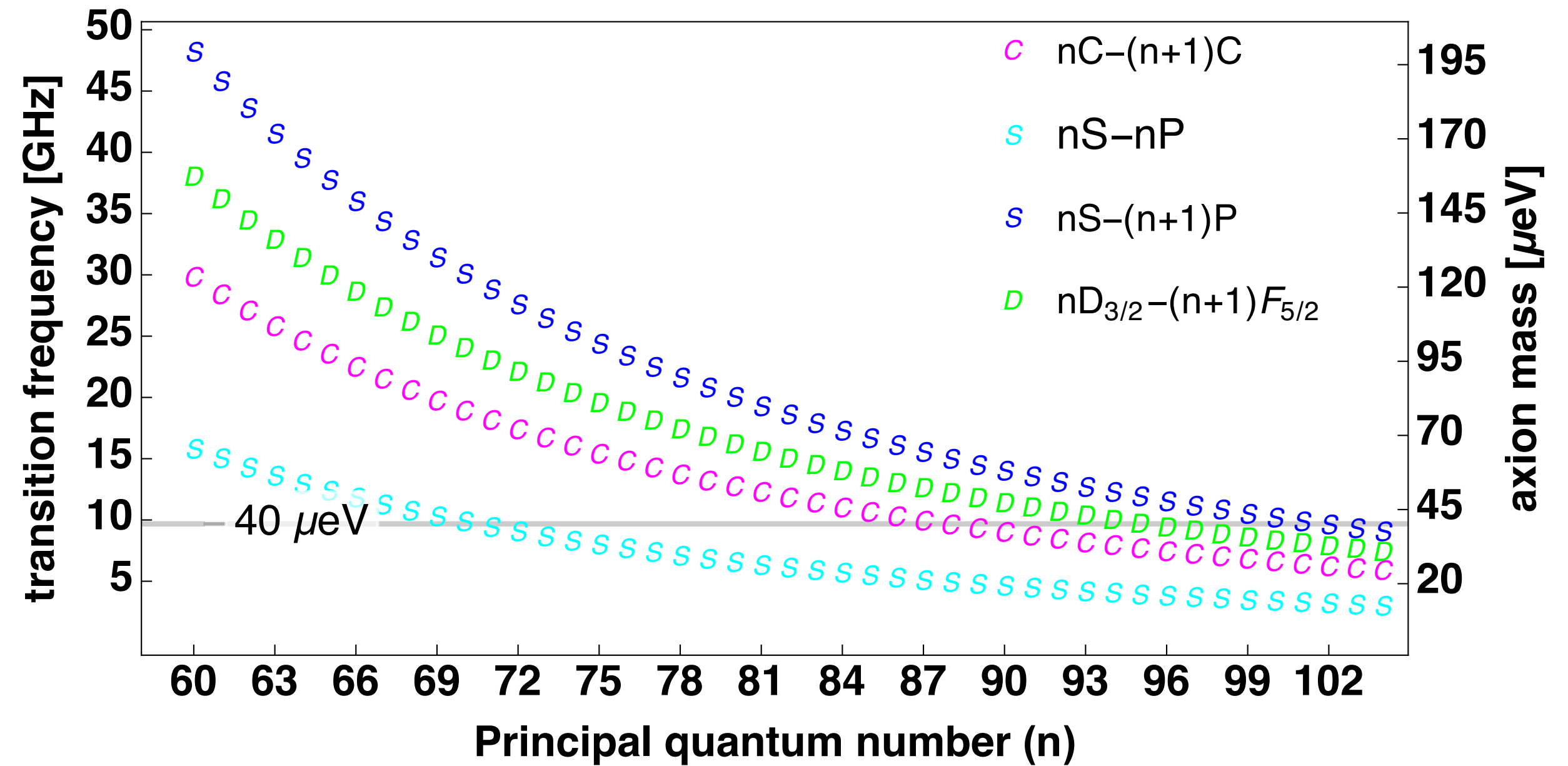
“Tuning” Rydberg atoms



Coarse-tuning: n

Fine-tuning: Zeeman/Stark

nC (circular, $l = |m| = n - 1$); nS ($l = 0$); nD ($l = 3$)

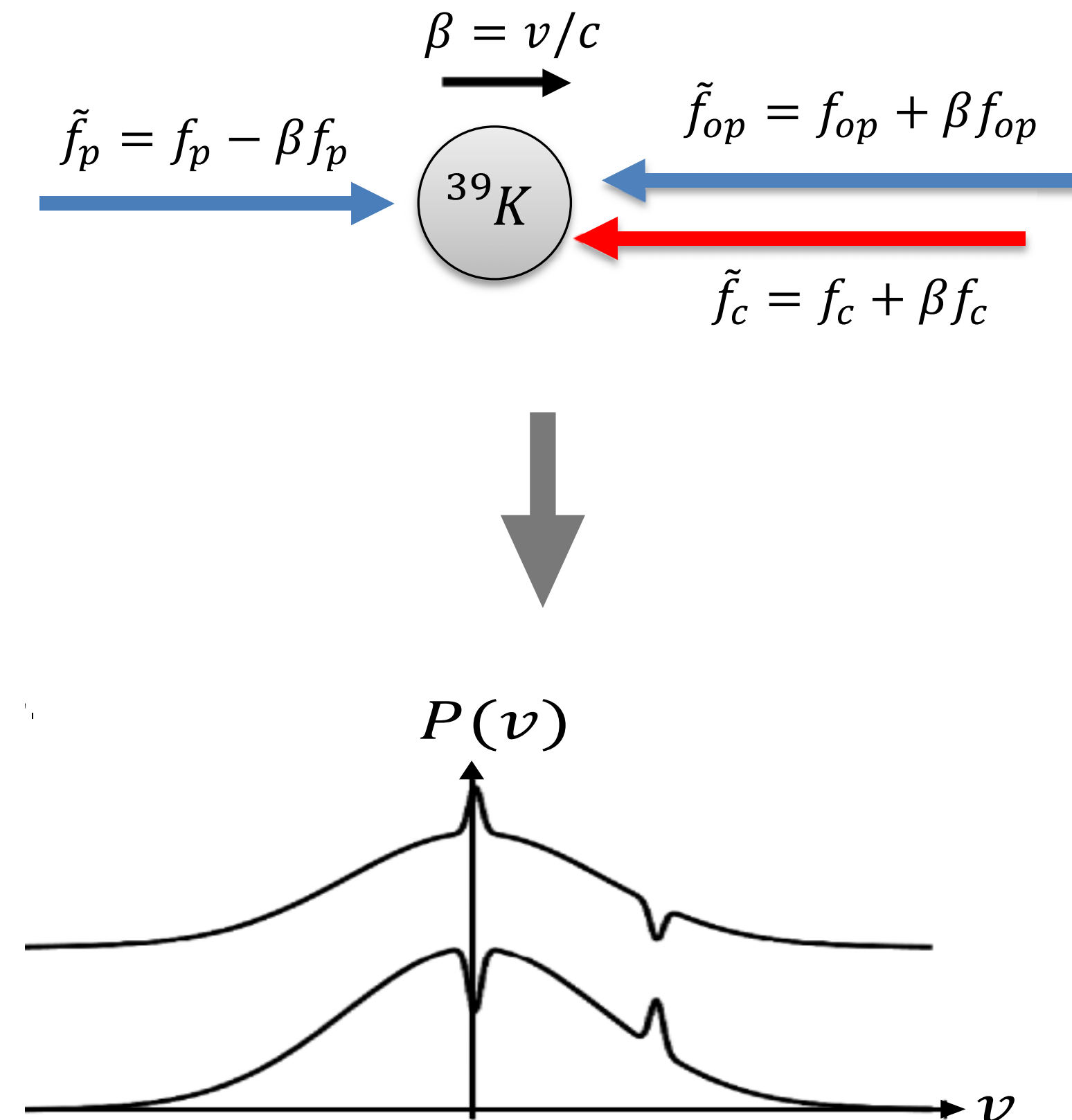
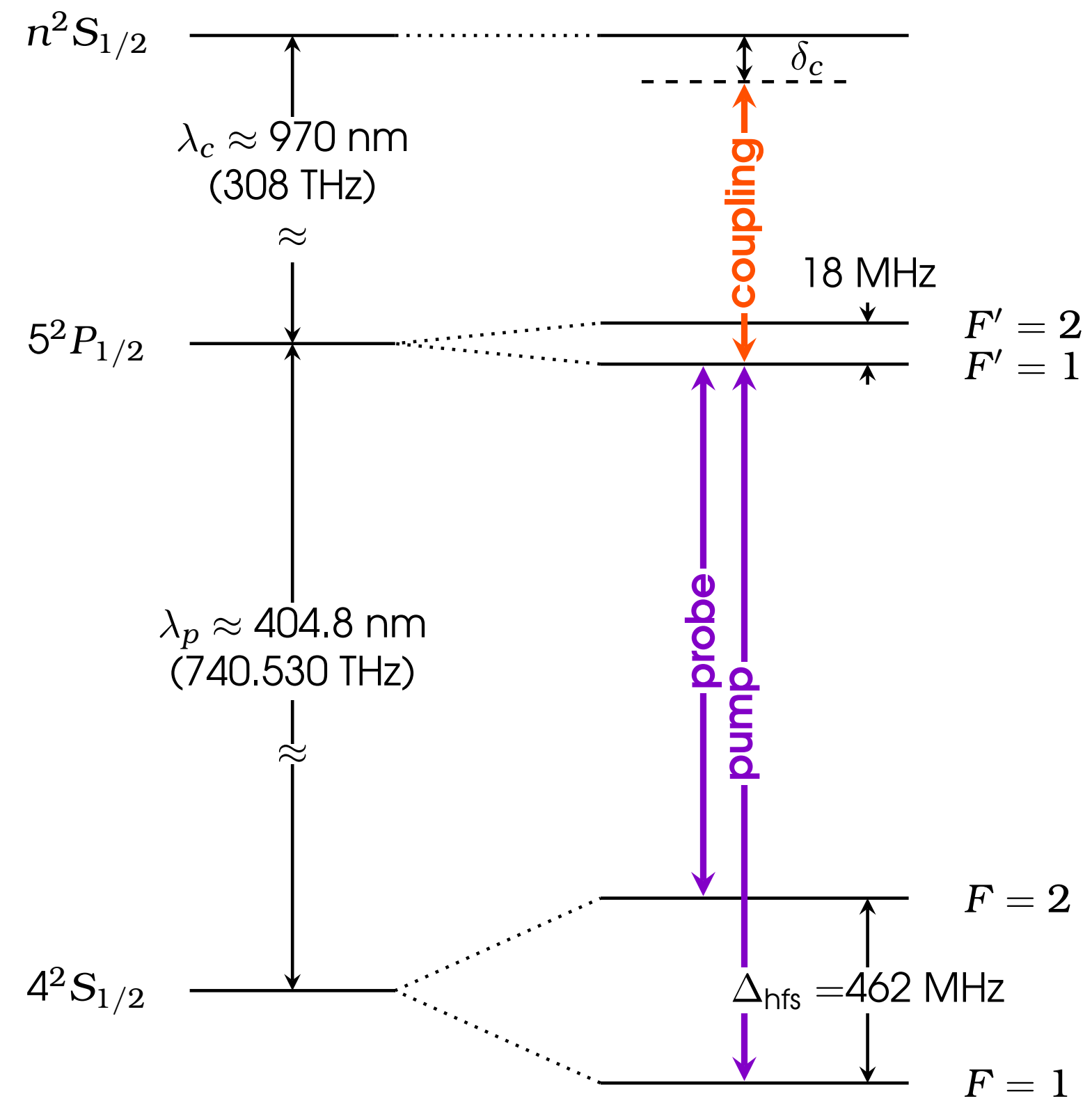


For $m_a = 40 \mu\text{eV} \approx 10 \text{ GHz}^*$:

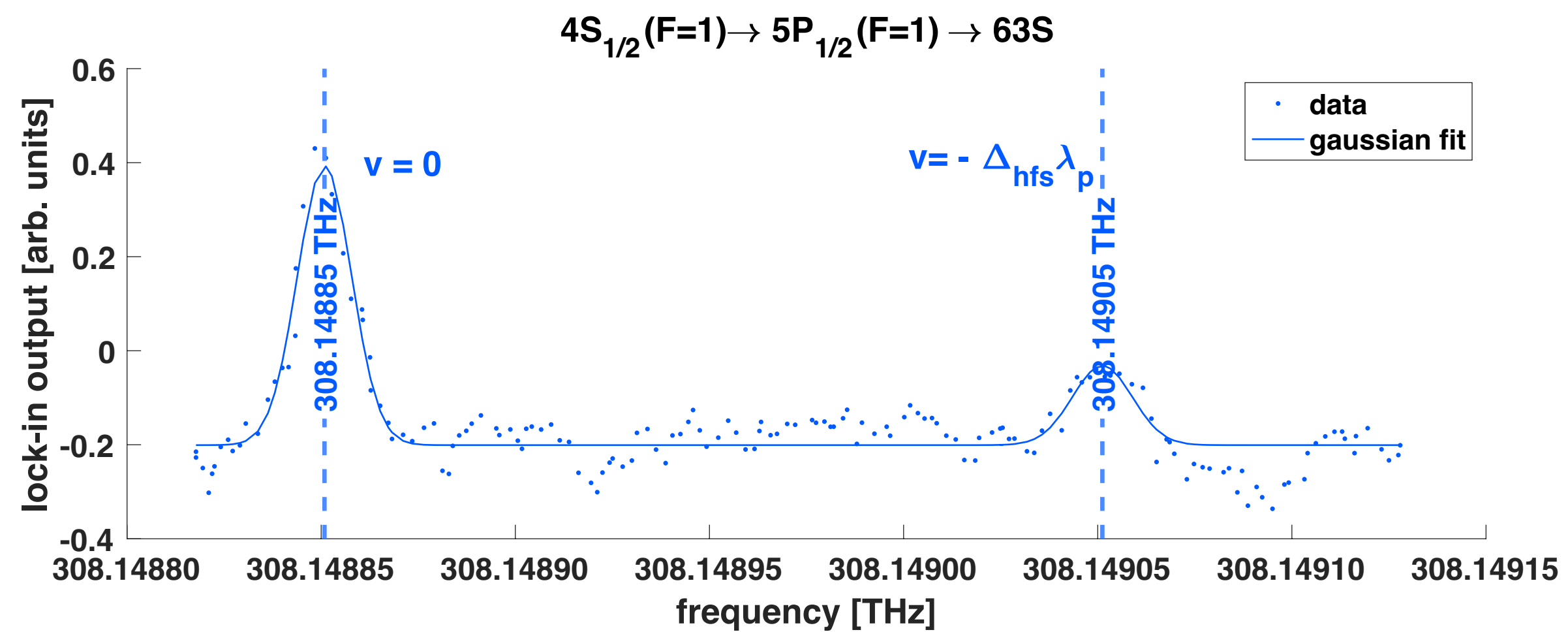
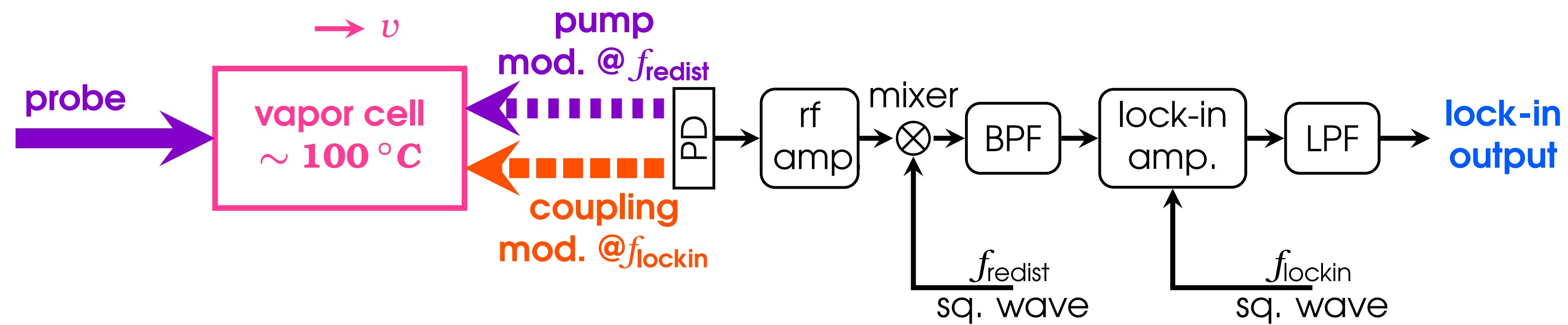
101S (and 70S), 95 $D_{3/2}$, and 87C

Rydberg spectroscopy

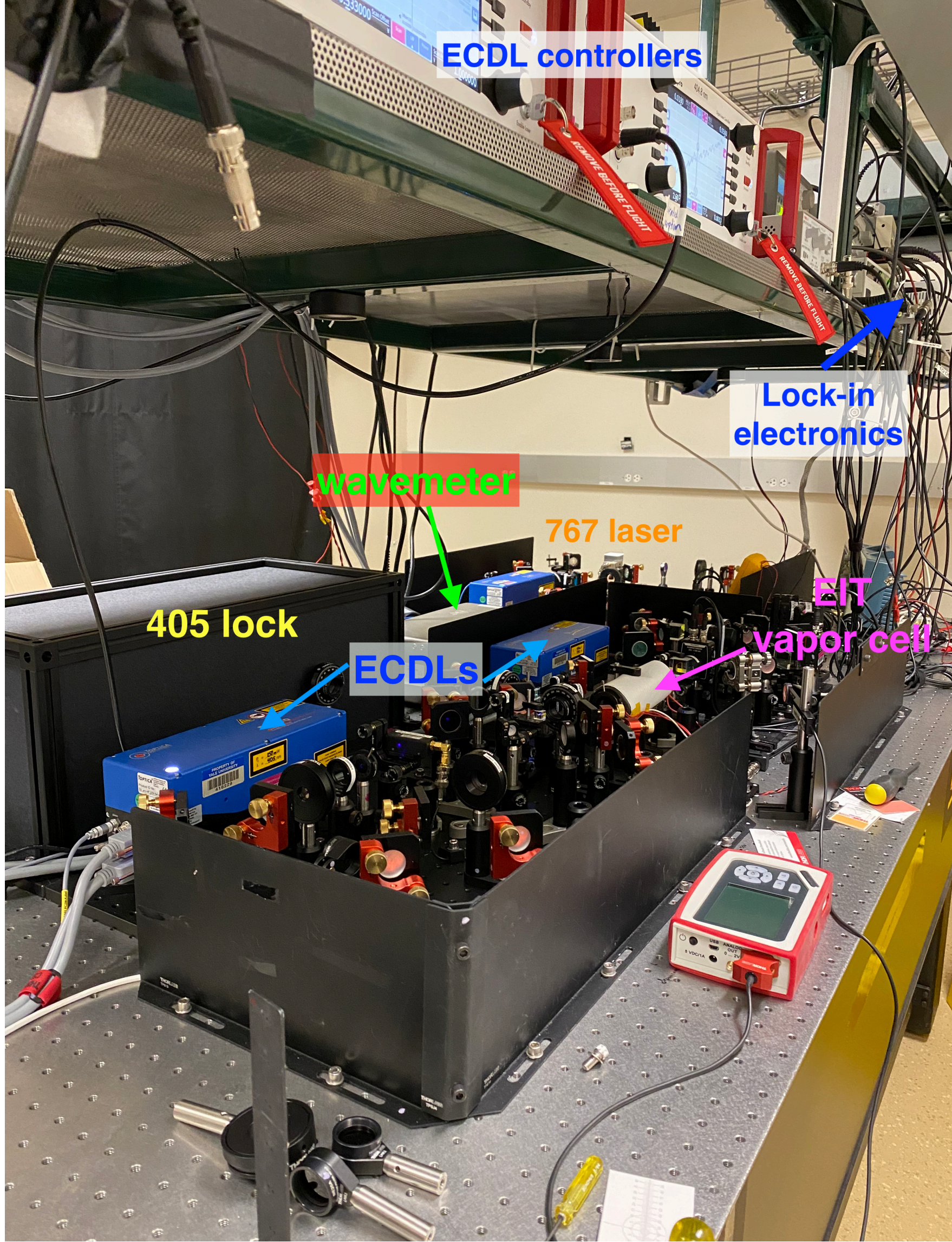
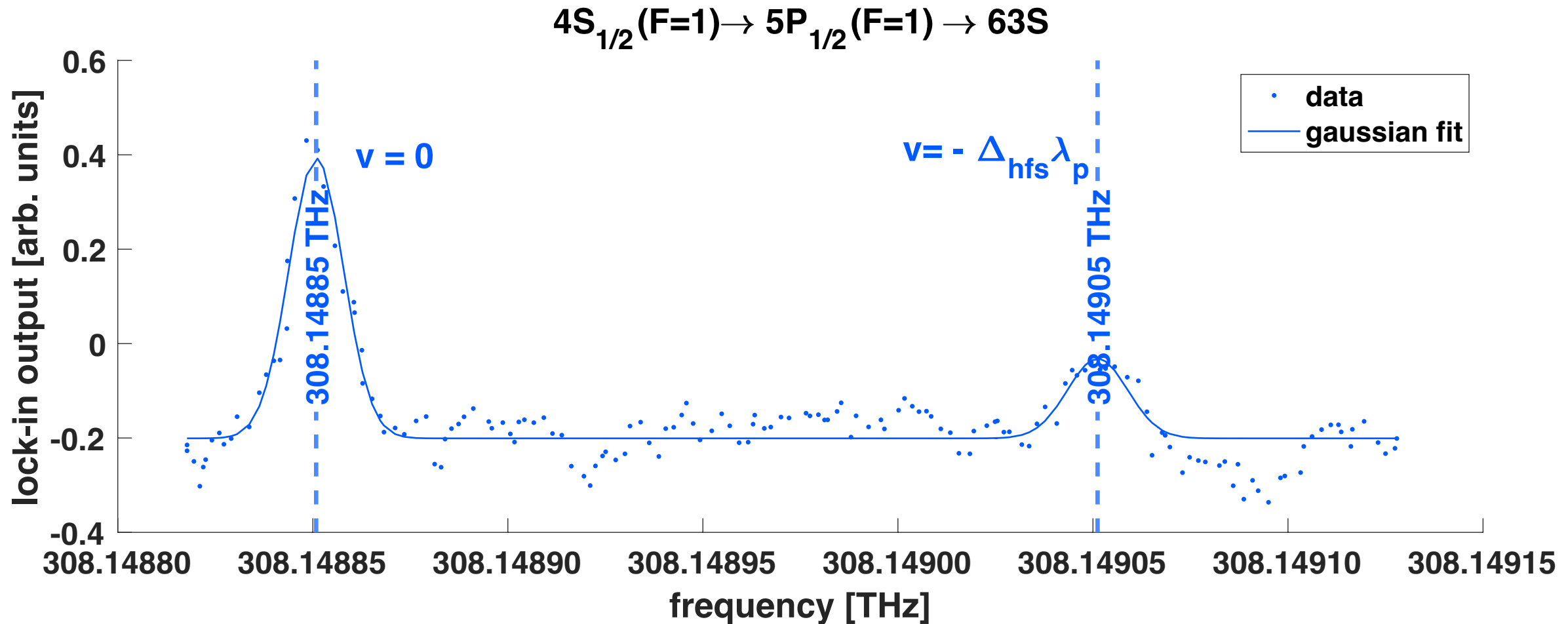
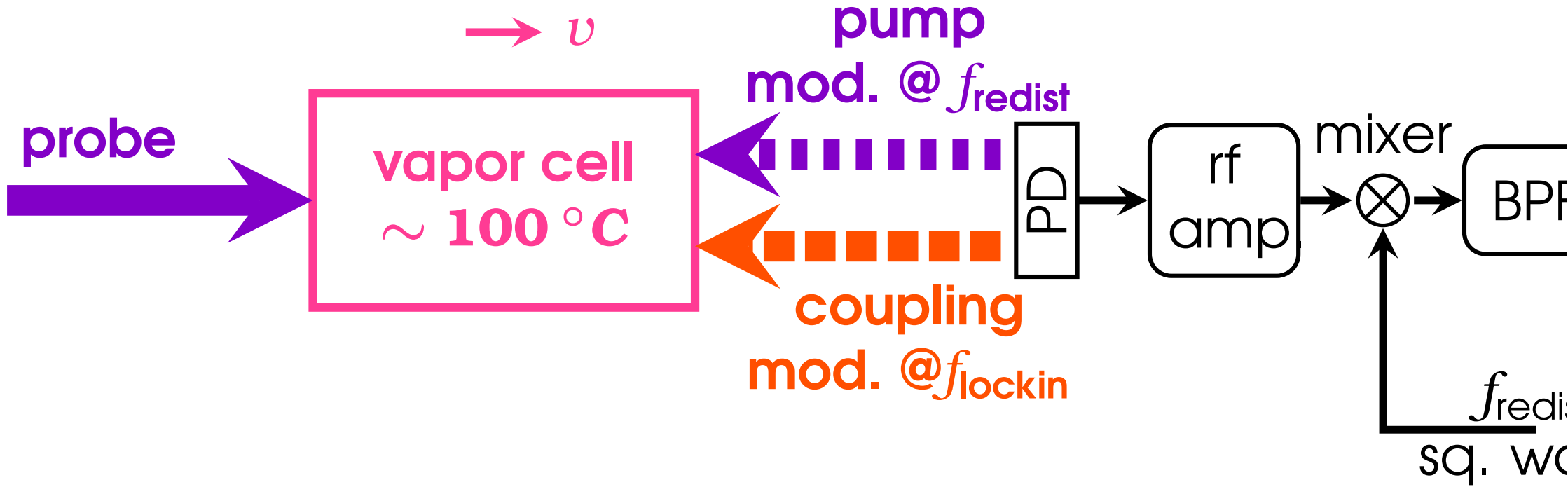
- **Goal:** identify the $n \sim 50 - 90$ transitions for 970 nm
- **Detection:** electromagnetically induced transparency (EIT)



EIT Setup

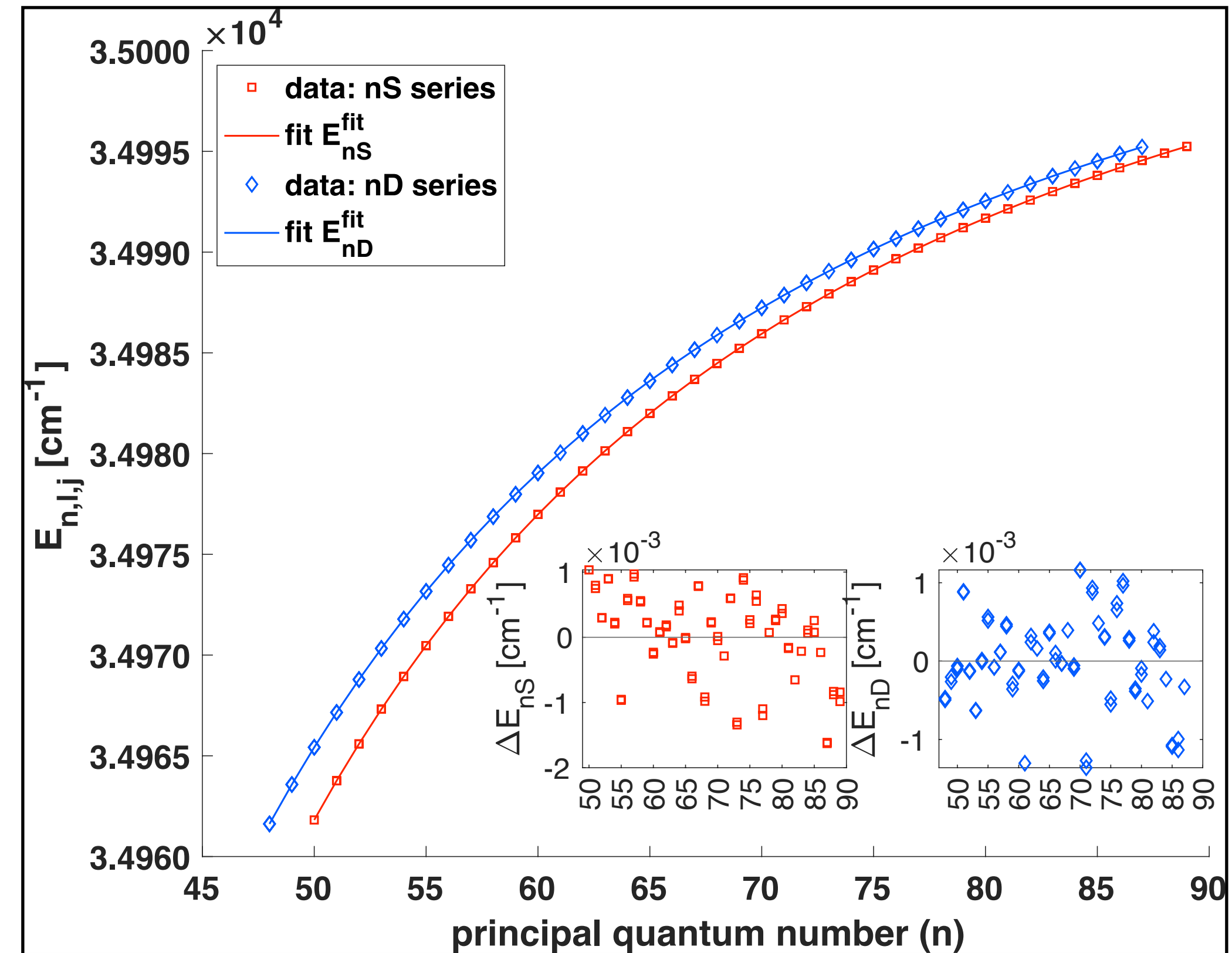


EIT Setup



$n = 47$ to 90 identified

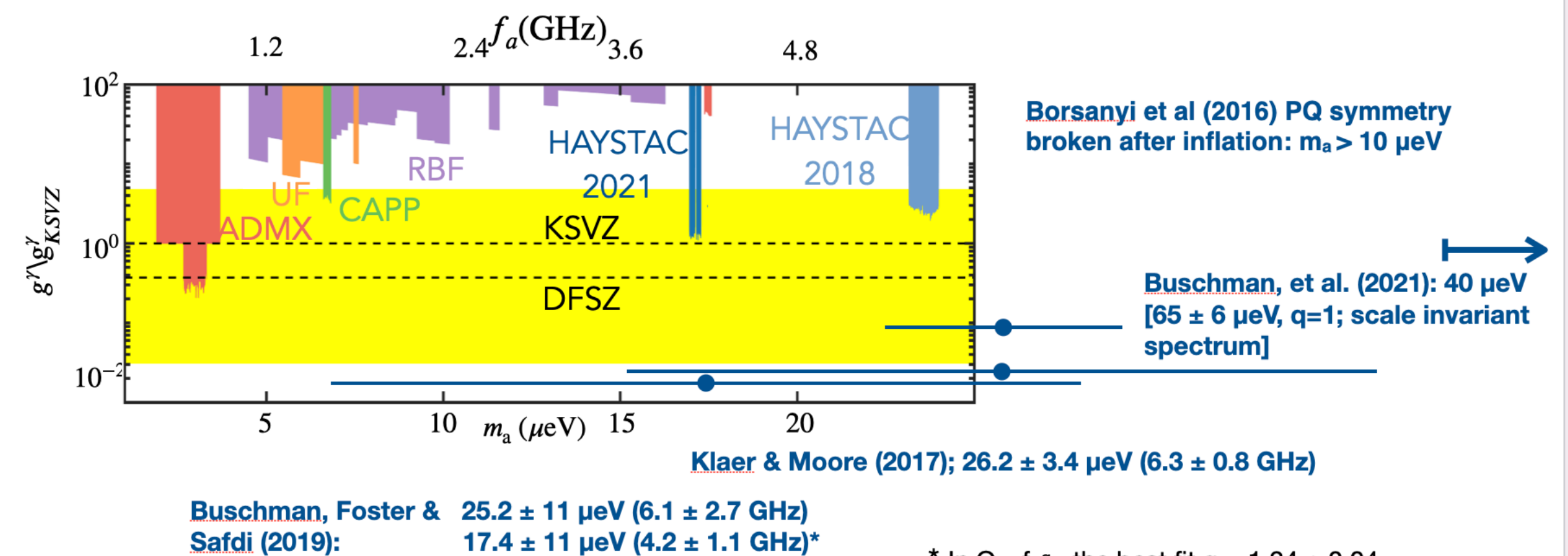
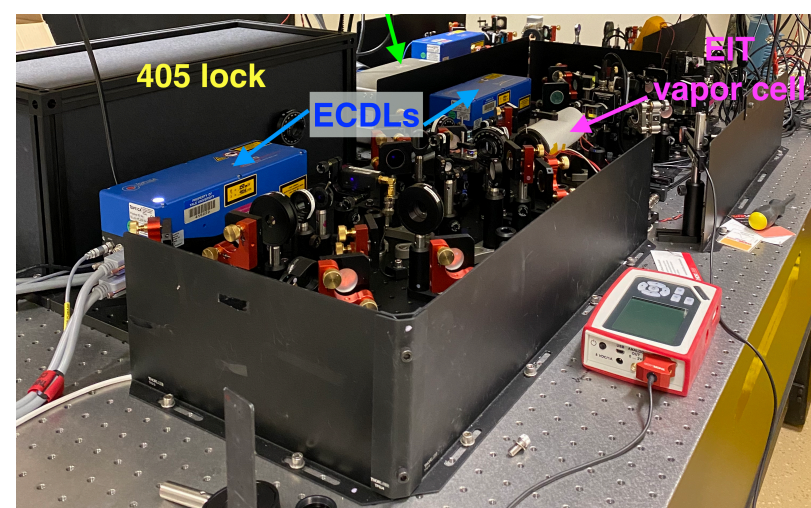
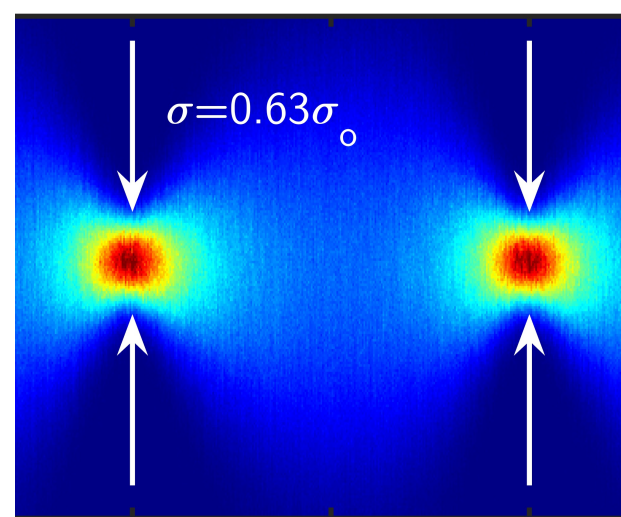
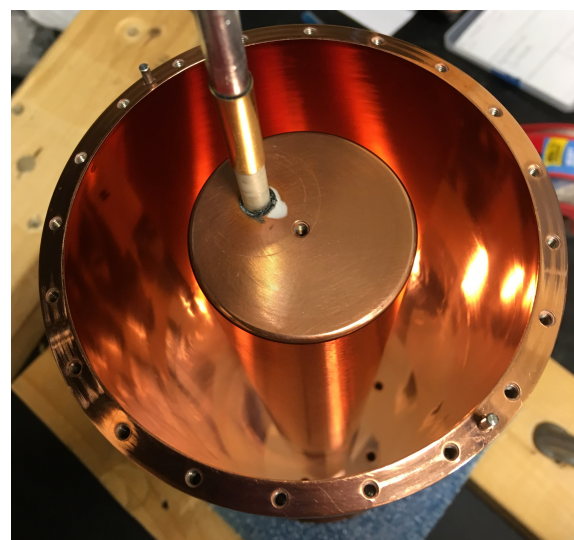
- ~ 80 transitions identified between $n = 47 - 90$
- 70 of the 80 newly observed Ryd. levels
- Enable access to $\sim 5 - 50$ GHz
($m_a \sim 20 - 200 \mu\text{eV}$)



YZ, S. Ghosh, S.B. Cahn, M.J. Jewell, D. H. Speller, RHM, arxiv:2112.04614

Conclusions

- HAYSTAC continues to scan 4 – 10 GHz
- Compelling case for axions at higher masses
- Rydberg atoms give us access to axion @ 10 – 1000 GHz/40 – 4000 μeV
- Effort @ Yale focus on 10 - 50 GHz
- $n \sim 50 - 90$ Rydberg levels observed
- Stay tuned



* In $\Omega_A \sim f_A^\alpha$, the best fit $\alpha = 1.24 \pm 0.04$ Rather than analytical 1.187